# SQUEEZING

# Field quadratures

$$\mathcal{E}_{x}(z,t) = \mathcal{E}_{0} \sin kz \sin(\omega t + \phi)$$
  
=  $\mathcal{E}_{0} \sin kz (\cos \phi \sin \omega t + \sin \phi \cos \omega t)$   
=  $\mathcal{E}_{1} \sin \omega t + \mathcal{E}_{2} \cos \omega t$   
 $\mathcal{E}_{1} = \mathcal{E}_{0} \sin kx \cos \phi; \quad \mathcal{E}_{2} = \mathcal{E}_{0} \sin kx \sin \phi$ 

$$X_{1}(t) = \sqrt{\frac{\varepsilon_{0}V}{4\hbar\omega}} \mathcal{E}_{0} \sin\omega t$$
$$X_{2}(t) = \sqrt{\frac{\varepsilon_{0}V}{4\hbar\omega}} \mathcal{E}_{0} \cos\omega t$$
$$\mathcal{E}_{x}(z,t) = \sqrt{\frac{4\hbar\omega}{\varepsilon_{0}V}} \sin kx (\cos\phi X_{1}(t) + \sin\phi X_{2}(t))$$

# Phasor representation



#### Reminder: Light waves as harmonic oscillators

New coordinates

$$q(t) = \sqrt{\frac{\varepsilon_0 V}{2\omega^2}} \mathcal{E}_0 \sin \omega t$$
$$p(t) = \sqrt{\frac{V}{2\mu_0}} B_0 \cos \omega t = \sqrt{\frac{\varepsilon_0 V}{2}} \mathcal{E}_0 \cos \omega t$$
$$p = \dot{q}$$
$$\ddot{q} = \dot{p} = -\omega^2 q$$
$$E_{em} = \frac{1}{2} \left( p^2 + \omega^2 q \right)$$

$$X_1(t) = \sqrt{\frac{\omega}{2\hbar}}q(t)$$
  $X_2(t) = \sqrt{\frac{1}{2\hbar\omega}}p(t)$ 

$$q(t) = \sqrt{mx(t)}$$
$$p(t) = \frac{1}{\sqrt{m}} p_x(t)$$

# Field Quadratures

$$X_1(t) = \sqrt{\frac{\omega}{2\hbar}}q(t)$$
  $X_2(t) = \sqrt{\frac{1}{2\hbar\omega}}p(t)$ 

$$\hat{q}(t) = \sqrt{\frac{\hbar}{2\omega}} (\hat{a} + \hat{a}^{\dagger})$$
$$\hat{p}(t) = \sqrt{\frac{1}{2\hbar\omega}} (\hat{a} - \hat{a}^{\dagger})$$

$$\hat{X}_1(t) = \frac{1}{2}(\hat{a} + \hat{a}^{\dagger}) \qquad \hat{X}_2(t) = \frac{1}{2}(\hat{a} - \hat{a}^{\dagger})$$

#### Uncertainty on field quadratures







Coherent states, shot noise and number-phase uncertainty

$$\alpha = X_{1} + iX_{2} = |\alpha|e^{i\phi} \quad (\text{with } |\alpha| = \sqrt{X_{1}^{2} + X_{2}^{2}})$$

$$\Delta X_{1} = \Delta X_{2} = \frac{1}{2}$$

$$|\alpha|^{2} = \overline{n}$$

$$\Delta n = \left(|\alpha| + \frac{1}{4}\right)^{2} - \left(|\alpha| - \frac{1}{4}\right)^{2} = |\alpha| = \sqrt{\overline{n}}$$

$$\Delta \phi = \frac{\text{uncertainty diameter}}{\alpha} = \frac{\frac{1}{2}}{\sqrt{\overline{n}}}$$

$$\Delta n\Delta \phi \ge \frac{1}{2}$$

$$\Delta n\Delta \phi \ge \frac{1}{2}$$

$$\Delta \overline{n} = \frac{\Delta \overline{n}}{X_{1}} \quad (\Delta \overline{n})$$

#### Laser Interferometer Gravitational Wave Observatory









![](_page_11_Figure_1.jpeg)

#### LISA

![](_page_12_Picture_1.jpeg)

http://lisa.nasa.gov

![](_page_13_Picture_1.jpeg)

![](_page_14_Figure_1.jpeg)

![](_page_15_Figure_1.jpeg)

![](_page_16_Figure_1.jpeg)

![](_page_17_Figure_0.jpeg)

## Vacuum

![](_page_18_Figure_1.jpeg)

#### Number states

![](_page_19_Figure_1.jpeg)

$$\begin{aligned} & \mathsf{Squeezing in classical harmonic oscillator} \\ F(x) &= -kx(1 - \varepsilon \sin 2\omega_0 t), \quad \omega_0 = \sqrt{\frac{k}{m}} \\ & \mathsf{Modulated force} \\ U(x) &= \frac{1}{2}kx^2(1 - \varepsilon \sin 2\omega_0 t) \\ & \mathsf{F}(x) = -\omega_0^2 x\varepsilon \sin 2\omega_0 t \\ & \mathsf{Equation of motion} \\ x(t) &= B(t)\cos\omega_0 t + C(t)\sin\omega_0 t \\ & \mathsf{Trial solution and its 2nd derivative} \\ & \ddot{x}(t) &= -\omega_0^2 x(t) - \omega_0 \dot{B}(t)\sin\omega_0 t + \omega_0 \dot{C}(t)\cos\omega_0 t \\ & \mathsf{B}(\mathsf{C} \ \mathrm{slowly \ varying, B''=C''=0}) \\ & \mathsf{Substitute into equation of motion} \\ & -\omega_0 \dot{B}(t)\sin\omega_0 t + \omega_0 \dot{C}(t)\cos\omega_0 t \\ &= -\omega_0^2 \varepsilon \sin 2\omega_0 t \Big[ B(t)\cos\omega_0 t + C(t)\sin\omega_0 t \Big] \\ & \mathsf{Apply \ trigonometry \ to \ find \ this} \\ & -\dot{B}(t)\sin\omega_0 t + \dot{C}(t)\cos\omega_0 t \\ &= -\frac{\omega_0 \varepsilon}{2} \Big[ B(t)\sin\omega_0 t + C(t)\cos\omega_0 t \Big] \\ & \dot{B}(t) \\ &= +\frac{\omega_0 \varepsilon}{2} B(t); \quad B(t) \\ & = B_0 e^{\frac{\varepsilon\omega_0 t}{2}} \\ & \mathsf{Equate \ sin \ and \ cos \ components} \\ & \dot{C}(t) \\ &= -\frac{\omega_0 \varepsilon}{2} C(t); \quad C(t) \\ & = C_0 e^{-\frac{\varepsilon\omega_0 t}{2}} \end{aligned}$$

#### Displacement operator and creation of coherent states

$$\hat{D}(\alpha) = e^{\alpha \hat{a}^{+} - \alpha^{*} \hat{a}}$$

 $\hat{D}^{-1}(\alpha)\hat{a}\hat{D}(\alpha) = \hat{a} + \alpha$ 

$$\begin{split} \hat{D}^{-1}(\alpha)\hat{a}\hat{D}(\alpha)\hat{D}^{-1}(\alpha)|\alpha\rangle &= \hat{D}^{-1}(\alpha)\hat{a}|\alpha\rangle = \alpha\hat{D}^{-1}(\alpha)|\alpha\rangle \\ &= (\hat{a}+\alpha)\hat{D}^{-1}(\alpha)|\alpha\rangle \end{split}$$

$$\hat{a}\hat{D}^{-1}(\alpha)|\alpha\rangle = 0$$
$$\hat{D}^{-1}(\alpha)|\alpha\rangle = |0\rangle$$
$$\hat{D}(\alpha)|0\rangle = |\alpha\rangle$$

### **Displacement of Vacuum**

![](_page_22_Figure_1.jpeg)

#### Squeezing operator

$$\hat{S}(z) = e^{\frac{1}{2}(z\hat{a}^2 - z^*\hat{a}^{+2})}$$
  $z = re^{i\varphi}$ 

$$\begin{aligned} \hat{t} &= \hat{S}(z)\hat{a}\hat{S}^{+}(z) = \hat{a} + z^{*}\hat{a}^{+} + \frac{1}{2!}|z|^{2}\hat{a} + \frac{1}{3!}|z|^{2}z^{*}\hat{a}^{+} + \frac{1}{4!}|z|^{4}\hat{a} + \dots \\ &= \hat{a}\left(1 + \frac{1}{2!}r^{2} + \frac{1}{4!}r^{4} + \dots\right) + \hat{a}^{+}e^{i\varphi}\left(r + \frac{1}{3!}r^{3} + \frac{1}{5!}r^{5} + \dots\right) \end{aligned}$$

 $\hat{t} = \hat{a}\cosh r + \hat{a}^{+}e^{i\varphi}\sinh r$ 

$$\hat{t}\hat{S}(z)|0\rangle = 0$$
  
$$\hat{t}\hat{S}(z)|\alpha\rangle = tS(z)|\alpha\rangle \qquad \qquad \hat{S}(z)|\alpha\rangle = \hat{S}(z)\hat{D}(\alpha)|0\rangle$$

#### **Displacement of Squeezed Vacuum**

![](_page_24_Figure_1.jpeg)

#### Squeezing operator

$$\hat{t} = \hat{a}\cosh r + \hat{a}^{+}e^{i\varphi}\sinh r$$

$$\hat{t}_{1} = \frac{1}{\sqrt{2}}(\hat{t}^{+} + \hat{t}) = \frac{1}{\sqrt{2}}(\hat{a}^{+} + \hat{a})(\cosh r + \sinh r) = \hat{X}_{1}e^{r}$$

$$\hat{t}_{2} = \frac{i}{\sqrt{2}}(\hat{t}^{+} - \hat{t}) = \frac{i}{\sqrt{2}}(\hat{a}^{+} - \hat{a})(\cosh r - \sinh r) = \hat{X}_{2}e^{-r}$$

$$\langle \alpha | \hat{S}^{+}(z)\hat{t}_{1}\hat{S}(z) | \alpha \rangle = \frac{\alpha}{2}e^{r}$$

$$\langle \alpha | \hat{S}^{+}(z)\hat{t}_{2}\hat{S}(z) | \alpha \rangle = \frac{\alpha}{2}e^{-r}$$

#### Generation of squeezed light

$$S(z) = e^{\frac{1}{2}(za^2 - z^*a^{+2})}$$

![](_page_26_Figure_2.jpeg)

![](_page_27_Figure_0.jpeg)

#### Detection of number squeezing

![](_page_28_Figure_1.jpeg)

#### Demonstrations of squeezed light I

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PHYSICAL REVIEW LETTERS

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#### Generation of Squeezed States by Parametric Down Conversion

Ling-An Wu, H. J. Kimble, J. L. Hall,<sup>(a)</sup> and Huifa Wu Department of Physics, University of Texas at Austin, Austin, Texas 78712

![](_page_29_Figure_6.jpeg)

![](_page_29_Figure_7.jpeg)

#### Demonstrations of squeezed light II

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#### **Pulsed Squeezed Light**

R. E. Slusher, P. Grangier, A. LaPorta, B. Yurke, and M. J. Potasek *AT&T Bell Laboratories, Murray Hill, New Jersey 07974* (Received 21 September 1987)

![](_page_30_Figure_6.jpeg)

![](_page_30_Figure_7.jpeg)

#### Quantum noise in amplifiers

![](_page_31_Figure_1.jpeg)

#### Exercises

- 1. Use the definition of field quadratures  $X_1(t)$  and  $X_2(t)$  to verify what are their physical dimensions.
- 2. A ruby laser operating at 693nm emits pulses of energy 1mJ. Calculate the uncertainty in the phase of the laser light .
- 3. The proposed Laser Interferometer Space Antenna (LISA) experiment for gravity wave detection will use a standard Michelson interferometer (i.e. no power reciclying or cavity enhancement) with a laser operating at 1064 nm The length if the arms of the interferometer is 5×10<sup>6</sup> km and the power of the beams that form the interference pattern is ~10<sup>-11</sup> W. Calculate the minimum strain that can be detected.
- 4. !! Demonstrate that an elastic force  $F = -kx(1 \varepsilon \sin 2\omega_0 t)$  modulated at twice the natural frequency ( $\varepsilon <<1$ ) produces on an object of mass *m* oscillations that grow in time at one quadrature, while are damped at the other quadrature.